

Catalogue of solar type II radio bursts observed from September 1990 to December 1993 and their statistical analysis^{*}

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Abstract. — Solar type II radio bursts represent the radio signature of shock waves travelling through the solar corona. They are associated with flares, coronal mass ejections (CME's) and interplanetary shocks. Type II radio bursts appear as emission stripes slowly drifting from high to low frequencies in dynamic radio spectra. The spectral features of all solar type II radio bursts observed by the new radiospectrograph of the Astrophysikalisches Institut Potsdam in Tremsdorf during the time period from September 1, 1990 to December 31, 1993, i.e., during the first part of the ULYSSES spacecraft mission, are summarized and statistically investigated.

Key words: Sun: activity; radio radiation

1. Introduction

Wild & McCready (1950) first reported the observations of so-called type II bursts occurring in nonthermal solar radio radiation. In dynamic radio spectra a type II burst appears as an emission band slowly drifting from high to low frequencies. An example of such a type II radio burst is presented in Fig. 1. These bursts are interpreted as the radio signature of shock waves in the solar corona. Such shock waves can be generated by flares and/or coronal mass ejections (CME's) (Roberts 1959; Smerd et al. 1975; Wild et al. 1972) (cf. also Dryer 1982; Nelson & Melrose 1985; Bougeret 1985; Aurass 1992; Mann 1995 as reviews). Sometimes coronal shock waves are able to penetrate into interplanetary space where they appear as interplanetary shocks. These interplanetary shocks may emit radio waves observed as interplanetary type II radio bursts observed below 1 MHz (Cane et al. 1982; Stone et al. 1984; Lengyel-Frey & Stone 1989).

Generally, solar type II radio bursts consist of two components, i.e., the so-called “backbone” and “herringbones” (cf. e.g. Krüger 1979). The “backbone” is the slowly drifting emission band. The drift rate is roughly -0.1 MHz/s (Nelson & Melrose 1985; Urbarz 1990). Most of the “backbones” show a fundamental-harmonic structure (Nelson & Melrose 1985). In rare cases the third harmonic band is also visible (cf. Fig. 1) (Roberts 1959; Kliem

et al. 1992; Aurass et al. 1995). Furthermore, the “backbone” has a great variety of fine structures, e.g., band splitting, multiple lane structures etc. (McLean 1967; Wild & Smerd 1972; Smerd et al. 1975; Nelson & Robinson 1975; Nelson & Melrose 1985). The “herringbones” are rapidly drifting emission stripes extending to high and low frequencies from the “backbone” in dynamic radio spectra. Twenty percent of all type II bursts are accompanied by “herringbones” (Roberts 1959). They resemble solar type III radio bursts and are interpreted as the radio signature of electron beams accelerated at the shock wave associated with the type II burst (Nelson & Melrose 1985). Cairns & Robinson (1987) deduced that these electrons have velocities in the range $0.08-0.4c$ (c , velocity of light).

Generally, it is assumed that the radio radiation is generated by suprathermal and/or high energy electrons. These electrons excite high frequency electrostatic plasma waves, i.e., upper hybrid waves or Langmuir waves, which scatter on ion density fluctuations and/or low frequency plasma waves and, thus, are converted into electromagnetic (radio) waves escaping the solar corona. Thus, the emission occurs slightly above the local electron plasma frequency. Furthermore, high frequency electrostatic waves can nonlinearly interact with themselves, coalescing into electromagnetic (radio) waves with frequencies approximately twice the electron plasma frequency. Thus, the former and latter mechanisms are responsible for the fundamental and harmonic radiation, respectively (cf. e.g. Krüger 1979). Benz & Thejappa (1988) argued

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Fig. 1. Solar type II radio burst observed in the range 40–800 MHz on September 27, 1993. It started at 12:10:00 UT. In this case the fundamental, first and second harmonic band of the “backbone” are clearly visible. The colour scale on the top panel indicates the intensity of the received radio radiation

that the radiation mechanism of the “backbone” is different from that of the “herringbones”.

Ginzburg & Zheleznyakov (1958) were the first to present a type II burst model based on electron streams accelerated at the shock front. Furthermore, Zheleznyakov (1965) and Zaitsev (1966) proposed that the electric current induced within the shock front becomes unstable and, subsequently, excites Langmuir waves needed for type II burst emission. Later on, Holman & Pesses (1983) and Benz & Thejappa (1988) suggested that solar type II radio bursts are generated by supercritical, quasi-perpendicular, fast magnetosonic shock waves in the corona. On the other hand, Mann & Lühr (1994) recently proposed a mechanism for producing solar type II burst radiation at supercritical, quasi-parallel shock waves.

In this paper the spectral properties of all solar type II radio bursts observed during the time period from September 1, 1990 until December 31, 1993 (cf. Table 1), i.e., since the launch of the ULYSSES satellite are reported. The observations were obtained with the new radiospectrograph of the Astrophysikalisches Institut Potsdam in Trensdorf. It consists of four sub-spectrometers working in the ranges 40–100 MHz, 100–170 MHz, 200–400 MHz and 400–800 MHz with a sweep rate of 10 Hz. Further details about this instrument are given by Mann et al. (1992). Comprehensive studies of observations of solar type II bursts and CME’s (Gosling et al. 1976; MacQueen 1980; Wagner 1980; Sheeley et al. 1984; Robinson et al. 1984) showed that CME’s with a velocity greater than 400 km/s, which is the case of 41% of all CME’s, are associated with solar type II radio bursts and interplanetary shocks, while

30% of all solar type II radio bursts are accompanied neither by CME's nor by interplanetary shocks. Furthermore, some interplanetary shocks are not associated with solar type II radio bursts. Maxwell et al. (1985) reported on a CME event accompanied by a type II radio burst. They showed that the shock wave associated with the type II burst was piston-driven by the CME. On the other hand, Gopalswamy & Kundu (1992) recently concluded by means of simultaneous measurements of type II radio bursts and CME's that solar type II radio bursts are generated by flare blast waves and not by CME's. A detailed knowledge of the properties of solar type II radio bursts and their associated shock waves is important for studying solar related events in the heliosphere as done by the ULYSSES satellite during the time interval under consideration. The URAP experiment aboard ULYSSES (Stone et al. 1992) is able to receive the radio radiation in the range 1 kHz – 1 MHz. According to the heliospheric density model (Bougeret et al. 1984) the 1 MHz plasma frequency level corresponds to a distance of 0.1 AU = 21 solar radii from the Sun. It should be noted, that the 40 MHz plasma frequency level, i.e., the lowest frequency of the observations presented, corresponds to a distance of 2 solar radii from the Sun (Mann et al. 1992).

In Sect. 2 we describe the data analysis of the observed solar type II radio bursts. The results of the statistical investigation of our data is given in Sect. 3. The paper is closed with some final remarks (Sect. 4).

2. Data analysis

Table 1 includes all solar type II radio bursts detected by the radiospectrograph of the Astrophysikalisches Institut Potsdam in Tremsdorf in the range 40–800 MHz during the time period September 1, 1990 - December 31, 1993. Sixty five of the 100 solar type II radio bursts detected are analyzed to obtain their properties in the dynamic radio spectra, e.g., instantaneous bandwidth, drift rate etc.. These 65 type II bursts (cf. Table 2) are chosen because their spectral features are clearly visible and, consequently, their spectral properties can be determined from the observations.

Figure 2 shows a schematic sketch of the dynamic radio spectra of a solar type II radio burst. The instrument employed (Mann et al. 1992) records the radio intensity in terms of steps of a grey scale on a black-white film. The radio intensity is measured in arbitrary units. In the frequency range considered (40–800 MHz) the background is largely nonuniform due to the instrumental noise, the galactic background, and the quiet Sun. The gauging of the instrument showed that a jump of the intensity of roughly 3 dB can just yet seen as a jump from one grey step to the next one. The measurements of bursts reported here are carried out as follows. The time is divided in N equidistant intervals of 10 s, i.e., $t_{i+1} - t_i = 10$ s, in each individual dynamic radio spectrum containing a type II

radio burst. At each time t_i the frequencies $f_{l,i}^F$, $f_{u,i}^F$, $f_{l,i}^H$, and $f_{u,i}^H$ are determined from the measurements (cf. Fig. 2) in such a way that the intensity is increasing above the background level by one grey scale step, i.e., an increase of approximately 3 dB, at the threshold. In some cases of the type II bursts investigated the peak intensity exceeds the saturation level, so that the frequency at the peak intensity can not be determined. All parameters are derived from this data set.

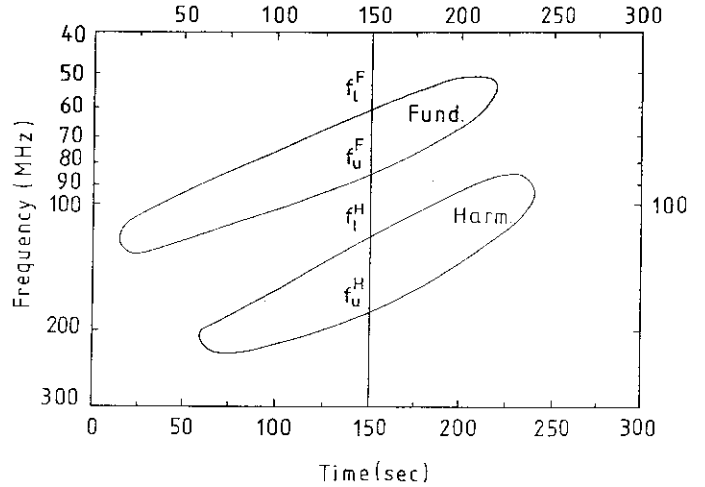


Fig. 2. Schematic sketch of the dynamic radio spectrum of a solar type II radio burst

The mean instantaneous width of the fundamental and harmonic band, i.e., $\langle(\Delta f/f)_F\rangle$ and $\langle(\Delta f/f)_H\rangle$, are calculated by

$$\left\langle\left(\frac{\Delta f}{f}\right)_F\right\rangle = \frac{1}{N} \sum_{i=1}^N \frac{f_{u,i}^F - f_{l,i}^F}{f_{l,i}^F}$$

and

$$\left\langle\left(\frac{\Delta f}{f}\right)_H\right\rangle = \frac{1}{N} \sum_{i=1}^N \frac{f_{u,i}^H - f_{l,i}^H}{f_{l,i}^H} \quad (1)$$

respectively. The linear regression analysis of the pairs $\{f_{l,i}^F, t_i\}$, $\{f_{u,i}^F, t_i\}$, $\{f_{l,i}^H, t_i\}$ and $\{f_{u,i}^H, t_i\}$ for all i 's ($i = 1, \dots, N$) provides the linear functions $f_k^F = a_k^F + b_k^F t$ and $f_k^H = a_k^H + b_k^H t$ with $k = u, l$, respectively. Then, the drift rates of the fundamental and harmonic band, $D_{f,F}$ and $D_{f,H}$, are determined by

$$D_{f,F} = \frac{b_l^F + b_u^F}{2} \quad \text{and} \quad D_{f,H} = \frac{b_l^H + b_u^H}{4} \quad (2)$$

respectively. (On the other hand, the drift rate could also be derived from the frequencies at the peak intensity. But this method was not used since the peak intensity can not accurately be determined for all cases.) Generally, the

Table 1. List of all solar type II radio bursts observed by the radiospectrograph of the Astrophysikalisches Institut Potsdam in Tremsdorf in the time period September 1, 1990 - December 31, 1993

date	start	end	kind of event	freq. range	date	start	end	kind of event	freq. range
90.09.14	1340.0	1419.3	II, IIIG	40–90	91.10.30	0628E	0820U	II, IV	100–170
90.10.01	1319.9	1327.3	II, C, U, IIIB	120–170	91.12.26	0759.6U	0906.0U	II?, IV, IIIGG	40–400
90.10.10	0654.9	0658.8	II, IIIG	100–170	91.12.26	0932.0	1014.2	IIIGG, IV, II, U, P	40–800
90.10.15	1226.8	1239.0	DCIM, II, C, IIIB	50–800	91.12.27	1106.7E	1125.5	IV, II, HARM, IIIG	40–800
90.12.10	0749.0	0758U	II, IIIG	100–170	91.12.28	1226.4	1227.3	II?	250–400
90.12.14	1333.9E	1343.3	II, IIIG	60–450	91.12.30	1013U	1035.3	IIIG, IV, II, HARM, SPIKES	40–800
90.12.20	0722.8U	0743.7	II, IV, P, RS	40–130	92.01.09	1240.3	1245.3	II, IV	100–400
90.12.23	0947.5	1000.0	IV, IIIG, II	40–400	92.02.02	0857.6	0951U	IIIGG, IV, II, H, P	50–800
90.12.30	1315U	1347U	IIIG, IV, II	100–350	92.02.07	1144.4	1319U	IIIG, IV, U, II?	40–800
91.01.08	1415.3	1416.5	II?	120–170	92.02.27	0951.2	1045U	II, H, HARM, IV, P, IIIG, Z, F, UG	40–800
91.01.18	1403.1	1421.6	IV, II, HARM	40–800	92.03.09	1336.4	1354.5	II, H, IV, F, IIIGG, RS	120–350
91.02.01	1144.4	1157.0	IIIG, IV, II	40–90	92.03.11	0905.9	0906.4	II?, (UNCLF)	100–300
91.02.19	1230.8	1239.4	II?	50–90	92.04.30	0808.0	0809.4	II?	40–50
91.02.25	0812.3	0843.5U	IIIGG, RSG, IV, II	40–400	92.05.07	0643.6	0654.5U	II, H, C	40–90
91.03.02	0735.2	0736.7	UNCLF, II?	120–150	92.06.08	0828.5	1115U	II, HARM, IV, I, F, P, Z, RS, IIIGG	40–800
91.03.04	1359.5	1413.8	UGG, II, H, HARM, IV	50–800	92.06.09	0605E	0639U	IIIG, IV, II, I, P, U?	100–800
91.03.05	0940.9	0943.5	II	40–140	92.06.16	0925.1	0927.5	II?	70–130
91.03.07	0747.3	1500U	UGG, IV, II, HARM, IIIGG, P, I, DC	40–800	92.07.16	0727.0	0731U	(II, H)?	200–400
91.03.12	1237.5	1257.1	U, IIIGG, IV, II, H, HARM, RS	40–800	92.08.03	0831.8	0837.0U	II?	200–400
91.03.16	1047.9	1105.7	SPIKES, IIIGG, IV, II, C, Z	40–800	92.07.14	1016.1	1040.7	IIIGG, RS, IV, P, II, H, HARM	40–400
91.03.16	1239.1	1259.2	IIIGG, IV, P, II?, SPIKES	40–550	92.08.10	1127.8	1128.5	II?	300–500
91.03.22	0828.8	0850.3U	IIIGG, IV, II, H	40–400	92.09.03	0930.8	0932.2	II	100–150
91.03.22	1146.1	1203.5	IIIGG, V, II?, IV	40–400	92.09.11	0604.6	0615U	IIIGG, IV, II, H, HARM	50–500
91.03.23	1227.0U	1247.0	IV, P, Z, F, IIIGG, II?	40–800	92.10.01	1039.1	1103U	IIIGG, II, H, HARM	40–400
91.03.29	0641.8	0701.1	SPIKES, IIIGG, U, IV, II, H, HARM	40–800	92.10.06	0950.1	1007.5U	IIIG, II?	40–170
91.04.02	1001.6	1021.0	IIIG, RS, UG, II, H, HARM, IV	40–700	92.10.06	1043.9	1046.0	IIIG, V, II?	40–400
91.04.11	1113.4	1142.5	IIIGG, II, H, HARM, IV, P	70–800	92.10.20	0909.5	0928U	IIIGG, IV, II, H, HARM, RSG	40–800
91.04.20	0959U	1023.0U	IV, II, H, P	40–170	92.10.28	1007.9U	1155U	IIIGG, IV, II, H, HARM, RS, P	40–800
91.04.29	1408.5	1416.0	II	100–170	92.11.09	1114.6	1120.5	IIIGG, (IV, II)?	40–170
91.05.09	0717.0U	0753U	II, H	40–170	92.12.09	1117.8	1125.2	II, H, HARM, IIIGG	100–600
91.05.16	0645.8	0737U	IIIGG, U, IV, II	40–800	92.12.30	0818.8	0832U	IIIGG, IV, II, RS, SPIKES	40–300
91.05.30	0937.0	0940.7	IIIB, RS, II, H, HARM	150–800	92.12.30	1031.8	1040U	II, HARM, C	40–90
91.06.01	1502.9	1504.8	II?	200–400	93.01.01	1150.6	1158.0	IIIGG, IV, II, H, HARM, RS, SPIKES	40–500
91.06.02	1412U	1757U	IV, IIIG, II?, C, P, F, Z	100–800	93.01.01	1335.6	1353.0	IIIGG, RS, IV, II, H, U	40–400
91.06.07	0616.5	0623U	II	40–90	93.01.07	0828.1	1005U	(II, HARM)?, IV, DCIM, P	60–800
91.06.10	1356.6U	1411.4U	IIIGG, II, HARM	40–400	93.01.15	0923.2	0930.6U	IV, II, H, U?	100–500
91.06.15	0814.2	1146U	II, H, HARM, IV, UG, IIIGG, P	40–800	93.02.18	0933.4	1019.6	II?, IIIGG, I, RS, C	40–400
91.06.21	0836.7	0907.0	IIIGG, RS, IV, II, H, HARM	40–500	93.02.18	1048.3	1111U	IIIG, IV, II, H, HARM	40–400
91.06.21	1237U	1252U	II?	40–90	93.03.04	1213.8	1311.9	IIIGG, IV, II?, U	40–500
91.06.28	0800.7	0813.0	U, IIIGG, IV, II, H, HARM	40–300	93.03.19	1242.4	1321.0U	II?, IIIGG, RS	40–400
91.06.28	1428.9	1435.8	IIIGG, IV, II, HARM, RSG	40–600	93.03.29	1019.2	1046.4	IIIGG, RS, II?	40–700
91.07.05	1002.8	1012.6	IIIGG, RS, IV, II, H	40–400	93.03.29	1251.3	1258.1	IV, IIIGG, U, II?	80–700
91.07.10	1203U	1531U	IIIGG, IV, II, U, I, F, Z	40–400	93.04.30	1427	1432	II	40–70
91.07.11	1036.3	1042.0	II	40–70	93.05.25	1513.0	1516.3	II	40–90
91.07.17	0627E	0652U	IIIGG, IV, II, H, HARM, RS	40–800	93.05.26	1029.2	1033.8	II?, I, IIIGG, RS	40–600
91.07.22	0946.0	1456U	IIIG, IV, U, II, H, HARM, P, RS	40–800	93.06.07	1410.4	1529U	IIIGG, IV, II, H, HARM, Z	40–170
91.09.02	1000.9	1005.5	II	40–60	93.07.22	0852.2	0915.8	II, H, HARM	40–150
91.09.08	0912.8	1015	IIIG, V, II, H, HARM, IV	40–350	93.09.27	1207.1	1219U	IIIG, II, H, HARM, P, DCIM	40–800
91.09.24	0750.5	0923U	IV, H, U, IIIGG	60–800	93.12.27	0911.6	0918.1	IIIG, II, HARM	110–500
91.10.05	0757.7	0823U	IIIG, IV, II, HARM	40–800	93.12.28	1208.1	1218.5	IIIG, II, H, HARM	40–700

Abbreviations at times: E, early as given; U, later than given event remarks: I, II, III, IV, V; type I, II, III, IV, V burst; UNCLF, unclassified events; B, single burst; S, storm bursts; U, type U bursts; G small group of bursts (< 10); GG, large group of bursts (> 10); C any continuum underlying other events; F, fluctuating continuum; RS, reverse drift bursts; H, herringbones; P, pulsations; Z, zebra bursts; HARM, harmonic emission; SPIKES, spike bursts; DCIM, events in the decimetric range

originally measured drift rates of the harmonic bands are roughly twice of those of the fundamental bands. In order to compare both the originally measured drift rate of the harmonic band is divided by 2 as done in Eq. (2) as well as in Table 2 and Sect. 3. Finally, the mean value of the ratio $\langle f_H/f_F \rangle$ for the middle frequencies of the harmonic and the fundamental bands is found by

$$\left\langle \frac{f_H}{f_F} \right\rangle = \frac{1}{N} \sum_{i=1}^N \frac{f_{l,i}^H + f_{u,i}^H}{f_{l,i}^F + f_{u,i}^F} \quad (3)$$

In Table 2 the mean values and the standard deviation of the instantaneous bandwidths $\Delta f/f$, the drift rate

D_f , and the ratio f_H/f_F are listed for all type II bursts analyzed. It should be mentioned that the instantaneous bandwidth and the drift rate are separately determined for the fundamental and harmonic bands as given in Table 2.

The method of determining the frequencies $f_{l,i}^F$, $f_{u,i}^F$, $f_{l,i}^H$, and $f_{u,i}^H$ leads to some errors in contrast to the determination at the 3 dB increase of the intensity. But this error can be considered to be small since our method is roughly the same. Furthermore, the calculation of the instantaneous bandwidth (cf. Eq. (1)), the drift rate (cf. Eq. (2)) and the ratio f_H/f_F (cf. Eq. (3)) is done by a statistical method, so that the above mentioned errors are reduced by this averaging.

Table 2. List of the spectral characteristics (general spectral structure (fundamental (F), harmonic (H)), instantaneous bandwidth $\langle \Delta f/f \rangle$ and drift rate $\langle D_f \rangle$ of the fundamental (left) and harmonic (right) band, respectively, as well as ratio $\langle f_H/f_F \rangle$) of the solar type II radio bursts investigated (cf. Table 1). The mean values and the standard deviations are given in the last panel

no.	date no. of day	universal time	structure	$\langle \frac{\Delta f}{f} \rangle$	$\langle D_f \rangle$ [MHz/s]	freq. range [MHz]	$\langle \frac{f_H}{f_F} \rangle$
1	90.09.14 257	1341-1349	F IIIGG, II	0.27 ± 0.05	-0.063 ± 0.002	66-50	
2	90.10.01 274	1325-1328	F IIIB, II	0.29 ± 0.06	-0.048 ± 0.023	63-55	—
3	90.10.10 283	0657-0659	F? IIIG, II	0.19 ± 0.03	-0.23 ± 0.01	145-110	—
4	90.10.15 288	1229-1231	F? II, IIIB	0.27 ± 0.05	-0.25 ± 0.02	143-103	—
5	90.12.10 344	0749-0803	F, H IIIG, II, IIIG	0.53 ± 0.13 — 0.53	-0.10 ± 0.04 — -0.10	81-45	2.03 ± 0.07 — 2.03
6	90.12.14 348	1335-1343	F, H IIIG, II	0.30 ± 0.12 0.20 0.43	-0.28 ± 0.11 -0.34 -0.22	125-110	1.94 ± 0.07 — 1.94
7	90.12.20 354	0725-0730	F, H RS, II, IV	0.16 ± 0.04 0.14 0.18	-0.07 ± 0.04 -0.08 -0.07	53-41	1.99 ± 0.13 1.87 2.11
8	90.12.23 357	0949-0957	F, H IV, IIIGG, II	0.36 ± 0.16 0.38 0.36	-0.18 ± 0.05 -0.22 -0.15	122-43	1.97 ± 0.23 1.85 2.06
9	90.12.30 364	1321-1324	H IIIG, IV, II	0.35 ± 0.02 — 0.35	-0.23 ± 0.06 — -0.23	≈ 140	—
10	91.01.18 18	1411-1421	F? VI, II	0.12 ± 0.04	-0.19 ± 0.07	156-52	—
11	91.02.01 32	1451-1456	F IIIG, IV, II	0.13 ± 0.04	-0.11 ± 0.01	82-55	—
12	91.02.25 56	0834-0838	F IIIGG, IV, II	0.30 ± 0.05	-0.07 ± 0.03	60-42	—
13	91.03.04 63	1402-1405	F, H UGG, IV, II	0.48 ± 0.39 0.20 0.75	-0.26 ± 0.02 -0.28 -0.25	133-103	1.81 ± 0.40 2.09 1.52
14	91.03.05 64	0941-0943	F, H II	0.16 ± 0.05 — 0.16	-0.10 ± 0.02 — -0.10	$\approx 90-60$	—
15	91.03.07 66	0756-0758	F, H IV, II, IIIGG	0.23 ± 0.08 0.17 0.28	-0.09 ± 0.08 -0.08 -0.10	53-40	2.06 ± 0.05 2.10 2.03
16	91.03.12 71	1252-1257	F, H IIIGG, IV, II	0.39 ± 0.10 0.39 0.31	-0.08 ± 0.01 -0.10 -0.06	62-46	2.08 ± 0.08 2.05 2.12
17	91.03.22 81	0844-0852	F? IIIGG, IV, II	0.25 ± 0.12 0.25 —	-0.13 ± 0.03 -0.13 —	118-63	—
18	91.03.23 82	1237-1239	H? IV, IIIGG	0.11 ± 0.03 — 0.11	-0.19 ± 0.03 — -0.19	160-150	—
19	91.03.29 88	0649-0657	F, H IV, II, IIIGG	0.52 ± 0.26 0.26 0.78	-0.12 ± 0.02 -0.10 -0.14	$\approx 100-49$	1.80 ± 0.59 2.22 1.38
20	91.04.02 92	1007-1009	F, H IIIG, II, IV	0.47 ± 0.20 0.33 0.61	-0.14 ± 0.03 -0.17 -0.11	68-50	2.10 ± 0.07 2.10 —
21	91.04.11 101	1116-1118	F, H IIIGG, II, IV	0.32 ± 0.21 0.17 0.46	-0.36 ± 0.02 -0.37 -0.34	149-101	1.82 ± 0.30 2.03 1.61
22	91.04.20 110	1007-1008	F? IV, II	—	-0.33 ± 0.10	80-43	—
				fund. harm.	fund. harm.		upper lower

Table 2. continued

no.	date no. of day	universal time	structure	$\left\langle \frac{\Delta f}{f} \right\rangle$	$\langle D_f \rangle$ [MHz/s]	freq. range [MHz]	$\left\langle \frac{f_H}{f_F} \right\rangle$
23	91.04.29 119	1409–1411	F? II	0.29 ± 0.04	-0.12 ± 0.06	120–107	—
24	91.05.09 129	0717–0721	F? II	0.20 ± 0.03	-0.08 ± 0.01	129–108	—
25	91.05.16 136	0708–0712	F IIIGG, IV, II	0.18 ± 0.07	-0.01 ± 0.004	55–49	—
26	91.05.30 150	0939–0941	F, H IIIB, II	0.31 ± 0.10 — 0.31	-0.48 ± 0.07 -0.41 -0.55	225–198	2.01 ± 0.06 2.01 —
27	91.06.02 153	1416–1418	F? IV, IIIG, II	0.16 ± 0.02	-0.13 ± 0.01	123–109	—
28	91.06.07 158	0617–0618	F? II	0.59 ± 0.07	-0.09 ± 0.05	52–45	—
29	91.06.10 161	1401–1407	F, H IIIGG, II	0.38 ± 0.10 0.36 0.43	-0.08 ± 0.02 -0.10 -0.06	74–47	2.07 ± 0.10 2.08 2.07
30	91.06.15 166	0816–0821	F, H II, IV, IIIGG	—	-0.37 ± 0.03 -0.35 -0.39	83–41	—
31	91.06.21 172	1237–1241	F II	0.12 ± 0.03	-0.04 ± 0.02	55–49	—
32	91.06.21 172	0849–0851	F IIIGG, IV, II	0.30 ± 0.12	-0.11 ± 0.08	55–42	—
33	91.06.28 179	0809–0811	F, H IIIGG, IV, II	0.45 ± 0.10 0.45 —	-0.13 ± 0.01 -0.13 -0.12	66–40	—
34	91.07.05 186	1008–1009	F? IIIGG, IV, II	0.25 ± 0.03	-0.15 ± 0.01	121–108	—
35	91.07.10 191	1218–1221	F? IIIGG, IV, II	0.07 ± 0.02	-0.14 ± 0.01	137–107	—
36	91.07.11 192	1036–1042	F II	0.27 ± 0.08	-0.04 ± 0.02	54–46	—
37	91.07.17 198	0639–0651	H IIIGG, IV, II	0.33 ± 0.05	-0.022 ± 0.001	40	—
38	91.07.22 203	1005–1014	F, H IIIG, IV, U, II	0.33 ± 0.04 0.33 0.33	-0.06 ± 0.06 -0.06 -0.07	63–44	2.16 ± 0.08 — 2.16
39	91.09.02 245	1001–1005	F II	0.17 ± 0.05	-0.04 ± 0.02	49–42	—
40	91.09.08 251	0917–0927	F, H IIIG, V, II, IV	0.38 ± 0.07 0.33 0.42	-0.10 ± 0.04 -0.12 -0.08	80–50	2.05 ± 0.21 1.90 2.13
41	91.09.24 267	0752–0800	H? IV, II, IIIGG	0.41 ± 0.08	-0.18 ± 0.1	133–52	—
42	91.12.26 360	0949–0955	F IIIGG, IV, II	0.32 ± 0.04	-0.049 ± 0.004	67–54	—
43	91.12.27 361	1114–1117	F, H IV, II, IIIG	0.31 ± 0.06 — 0.31	-0.05 ± 0.03 -0.05 -0.05	58–51	2.09 ± 0.06 2.05 2.13
44	91.12.30 364	1022–1031	F? IIIG, IV, II	—	-0.09 ± 0.03	81–51	—
				fund. harm.	fund. harm.		upper lower

Table 2. continued

no.	date no. of day	universal time	structure	$\left\langle \frac{\Delta f}{f} \right\rangle$	$\langle D_f \rangle$ [MHz/s]	freq. range [MHz]	$\left\langle \frac{f_H}{f_F} \right\rangle$
45	92.02.02 33	0910–0915	H IIIGG, IV, II	—	-0.07 ± 0.03	≈ 50	—
46	92.02.27 58	0951–0954	F, H	0.48 ± 0.40 0.19 0.76	-0.32 ± 0.08 $-0.24 -0.39$	113–56	1.88 ± 0.09 1.94 1.81
47	92.05.07 128	0643–0651	F II	0.27 ± 0.02	-0.043 ± 0.022	58–44	—
48	92.07.14 196	1025–1029	F? IIG, IV, IIG	—	-0.27 ± 0.13	≈ 100	—
49	92.08.03 216	0832–0835	H? II	0.37 ± 0.05	-0.22 ± 0.06	155–125	—
50	92.09.03 247	0931–0932	F? II	—	-0.21 ± 0.05	113–110	—
51	92.09.11 255	0607–0610	F? IIIGG, IV, II	0.23 ± 0.05	-0.39 ± 0.08	160–120	—
52	92.10.01 275	1048–1052	F, H IIIGG, II	0.38 ± 0.10 0.39 0.30	-0.07 ± 0.04 $-0.08 -0.06$	50–42	2.17 ± 0.13 — 2.17
53	92.10.20 294	0915–0924	F, H IIIGG, IV, II	0.74 ± 0.16 — 0.74	-0.13 ± 0.06 $-0.18 -0.07$	47–41	2.20 ± 0.10 2.20 —
54	92.10.28 302	1011–1021	F, H IIIGG, IV, II	0.25 ± 0.02 0.23 0.26	-0.07 ± 0.02 $-0.09 -0.05$	54–42	2.06 ± 0.09 2.06 —
55	92.12.29 364	1118–1123	F, H IIIGG, II	—	-0.38 ± 0.03 $-0.40 -0.35$	154–106	1.83 ± 0.11 — 1.83
56	92.12.30 365	1034–1039	F IIG, IV, II	0.32 ± 0.08	-0.071 ± 0.004	64–49	—
57	93.01.01 1	1344–1348	F III, IV, II, RS	0.42 ± 0.17	-0.25 ± 0.01	128–63	—
58	93.01.15 15	0927–0930	F? IV, II	0.34 ± 0.06	-0.37 ± 0.06	165–115	—
59	93.04.30 120	1428–1432	F II	0.24 ± 0.04	-0.068 ± 0.004	54–45	—
60	93.05.25 145	1514–1517	F II	0.27 ± 0.02	-0.060 ± 0.005	53–48	—
61	93.06.07 158	1425–1435	F, H	0.49 ± 0.16 0.42 0.51	-0.06 ± 0.03 $-0.07 -0.05$	50–43	2.12 ± 0.07 — 2.12
62	93.07.22 203	0853–0903	F, H	0.31 ± 0.10 — 0.31	-0.05 ± 0.04 $-0.05 -0.05$	66–48	2.11 ± 0.10 — 2.11
63	93.09.27 270	1210–1218	F, H IIG, II, DC	0.38 ± 0.02 0.36 0.40	-0.23 ± 0.02 $-0.22 -0.25$	136–54	1.95 ± 0.11 2.02 1.87
64	93.12.27 361	0915–0918	F, H IIG, II	0.20 ± 0.06 0.16 0.24	-0.20 ± 0.05 $-0.16 -0.25$	134–114	1.98 ± 0.08 2.03 1.92
65	93.12.28 362	1212–1216	F, H IIG, II	0.67 ± 0.12 — 0.67	-0.38 ± 0.08 $-0.46 -0.30$	120–65	1.99 ± 0.11 2.07 1.91
	average			0.32 ± 0.14 0.26 0.41	-0.16 ± 0.11 $-0.16 -0.17$		2.01 ± 0.11 2.04 1.95
				fund. harm.	fund. harm.		upper lower

3. Statistical analysis

The instantaneous bandwidths (cf. Eq. (1)) of the fundamental and harmonic bands of the type II bursts lie in the interval $0.07 \leq \langle(\Delta f/f)_F\rangle \leq 0.59$ with a mean value of $\langle\langle(\Delta f/f)_F\rangle\rangle = 0.26$ and $0.07 \leq \langle(\Delta f/f)_H\rangle \leq 0.74$ with a mean value of $\langle\langle(\Delta f/f)_H\rangle\rangle = 0.41$, respectively. The corresponding histograms are presented in Figs. 3 and 4. Combining the instantaneous bandwidths of both the fundamental and harmonic bands, a mean value and a standard deviation

$$\langle\langle(\Delta f/f)\rangle\rangle = 0.32 \pm 0.14$$

is obtained with the distribution given in the histogram of Fig. 5.

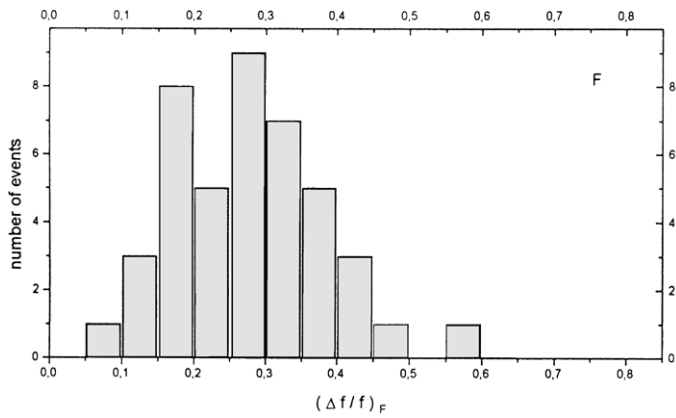


Fig. 3. Histogram of the mean instantaneous bandwidth $\langle(\Delta f/f)_F\rangle$ of the fundamental band of all solar type II radio bursts analyzed (cf. Table 2)

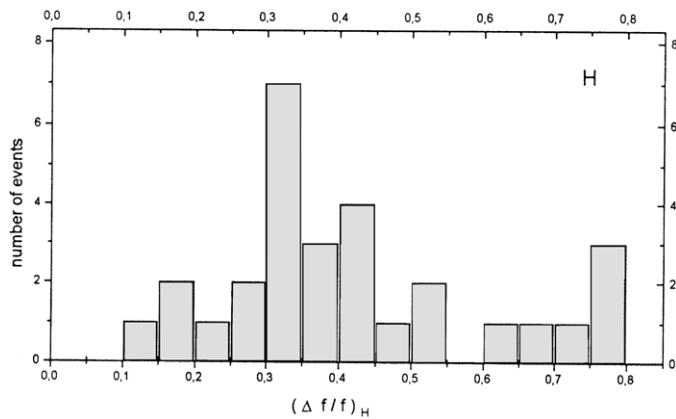


Fig. 4. Histogram of the mean instantaneous bandwidth $\langle(\Delta f/f)_H\rangle$ of the harmonic band of all solar type II radio bursts analyzed (cf. Table 2)

Generally, the mean value and standard deviation of the drift rate averaged over the whole set of solar type II

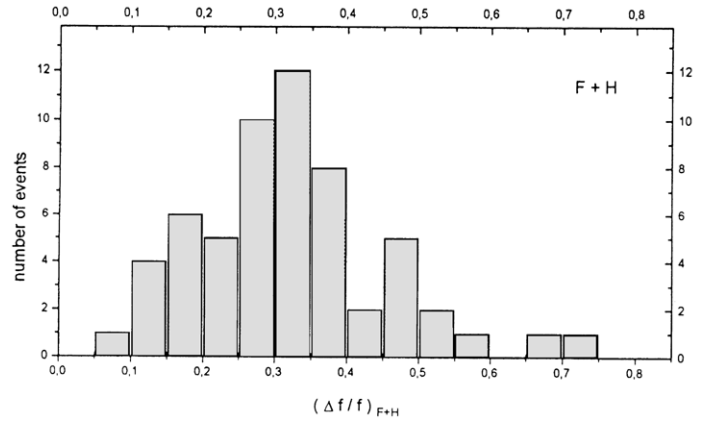


Fig. 5. Histogram of the instantaneous bandwidth of the fundamental and harmonic bands of all solar type II radio bursts analyzed (cf. Table 2)

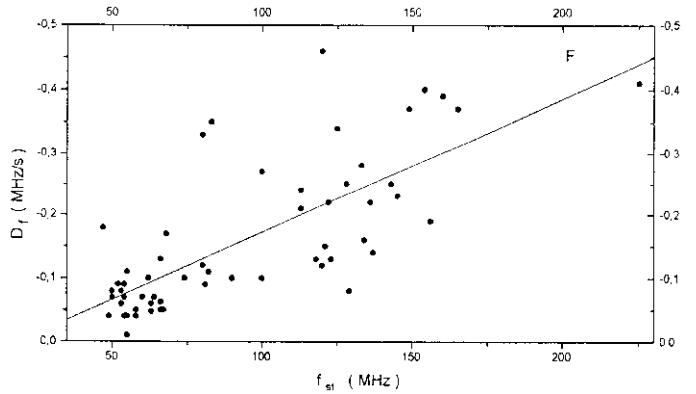


Fig. 6. Scatter plot of the magnitude of drift rate $-D_{f,F}$ versus the starting frequency f_{st} for the fundamental band of all individual solar type II radio bursts analyzed. The straight line results from a linear regression analysis with a correlation coefficient of 0.738

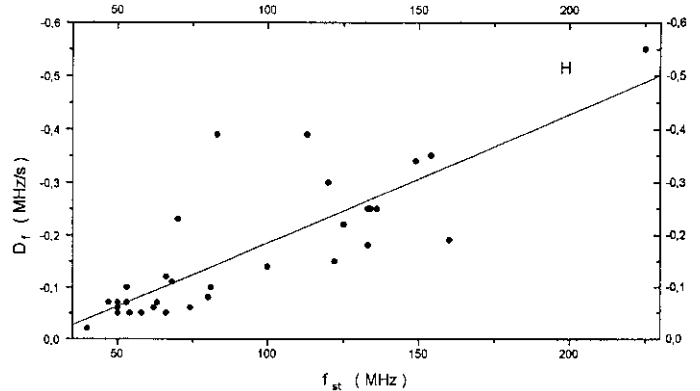


Fig. 7. Scatter plot of the magnitude of drift rate $-D_{f,H}$ versus the starting frequency f_{st} of the harmonic band of all individual solar type II radio bursts analyzed. The straight line results from a linear regression analysis with a correlation coefficient of 0.882

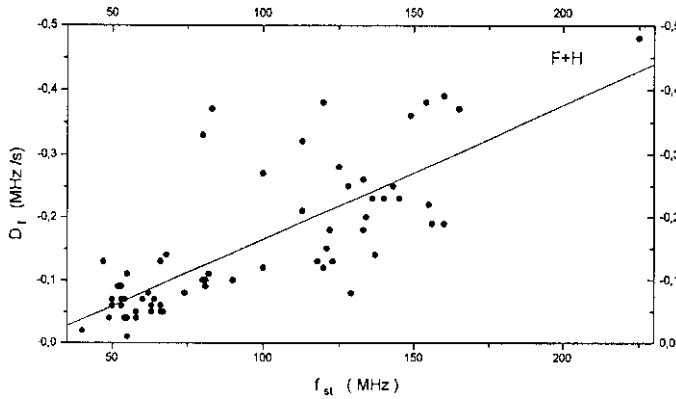


Fig. 8. Scatter plot of the magnitude of drift rate $-D_{f,F+H}$ versus the starting frequency f_{st} of both the fundamental and harmonic bands of all individual solar type II radio bursts analyzed. The straight line results from a linear regression analysis with a correlation coefficient of 0.779

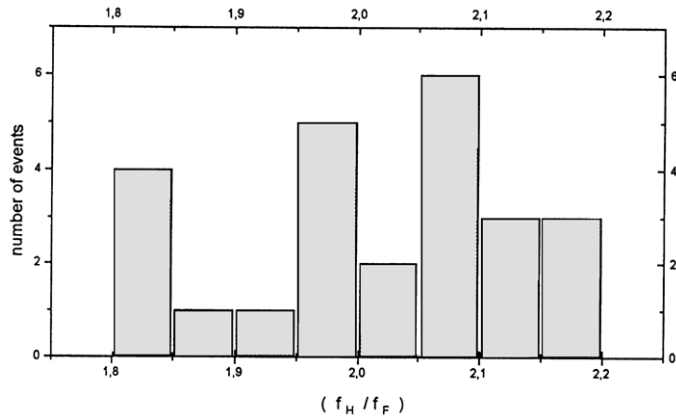


Fig. 9. Histogram of the mean values of the ratio $\langle f_H/f_F \rangle$ (cf. Eq. (3)) between the middle frequency of the harmonic and the fundamental band of the solar type II radio bursts analyzed (cf. Table 2)

radio bursts is found to be

$$\langle D_f \rangle = -0.16 \pm 0.11 \text{ MHz/s}$$

It has been shown that the magnitude of the drift rate increases with the starting frequency of type II bursts (cf. e.g. Urbarz 1990; Mann 1995). This property of type II bursts is confirmed by this study. The corresponding scatter plots of the magnitude of the drift rate $-D_f$ versus the starting frequency f_{st} of the individual bursts are presented for the fundamental and harmonic bands as well as for both data sets in the Figs. 6-8, respectively. Note that the half of the originally observed starting frequency has been used for f_{st} in the case of the harmonic bands. The linear regression analysis of the three data sets, i.e., $\{-D_{f,F}, f_{st}\}$, $\{-D_{f,H}, f_{st}\}$, $\{-D_{f,F+H}, f_{st}\}$ provides

$$(-D_{f,l}) = A_l + B_l f_{st} \quad (4)$$

($l = F, H, F+H$) with $A_F = -0.043 \text{ MHz/s}$, $A_H = -0.058 \text{ MHz/s}$, $A_{F+H} = -0.046 \text{ MHz/s}$, $B_F = 0.002 \text{ s}^{-1}$, $B_H = 0.002 \text{ s}^{-1}$, and $B_{F+H} = 0.002 \text{ s}^{-1}$ (cf. the straight lines in Figs. 6-8). Here, the data set $\{-D_{f,F+H}, f_{st}$ includes all drift rates of the fundamental and harmonic bands.

The analysis of the ratio $\langle f_H/f_F \rangle$ as defined by Eq. (3) provides the distribution presented in Fig. 9. The values of $\langle f_H/f_F \rangle$ lies in the interval $1.80 \leq \langle f_H/f_F \rangle \leq 2.20$ with a mean value and a standard deviation of

$$\langle \langle f_H/f_F \rangle \rangle = 2.01 \pm 0.11$$

As the Fig. 9 shows, the distribution is sharply localized around this mean value. This result slightly differs from those in previous papers (cf. e.g. Nelson & Melrose 1985). This difference is caused by the different methods of computing this ratio. In this paper the center-frequency of each emission band is used for calculating this ratio in contrast to the low-frequency edge of these bands as done in previous papers (cf. Maxwell & Thompson 1962). Taking into account this difference the result of this paper agrees well with that of Maxwell & Thompson (1962).

4. Final remarks

In this paper a catalogue of solar type II radio bursts observed during the time period September 1, 1990 until December 31, 1993 (cf. Table 1) is presented with several spectral properties (cf. Table 2) of each individual burst. A statistical analysis of these spectral properties has been performed (cf. Sect. 3). An interpretation of the statistical study of the instantaneous bandwidth has been given by Mann et al. (1995). They concluded that solar type II radio bursts are generated by supercritical, quasi-parallel, fast magnetosonic shocks travelling through the corona.

The observations were done during the first part of the ULYSSES mission. As mentioned above solar type II radio bursts are associated with CME's and interplanetary shocks (cf. Sect. 1), i.e., with large scale disturbances of the heliosphere. Such disturbances and associated particle events can be detected by the instruments aboard the ULYSSES satellite. Thus, the results presented in this paper will be useful for the investigation of solar related phenomena in the heliosphere by means of the data recorded by the ULYSSES satellite.

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