

# Stark broadening of solar Mg I lines<sup>\*</sup>

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**Abstract.** — Using a semiclassical approach, we have calculated electron-, proton-, Mg II-, Si II-, Fe II-, and Ar II- impact line widths and shifts for 267 Mg I multiplets, in order to provide the needed Stark broadening parameters for all important perturbers for investigation of Solar and laboratory plasma. An analysis of Solar Rydberg lines in the far infrared spectrum has been performed as well.

**Key words:** lines: profile-atomic and molecular data — Sun: photosphere — Sun: chromosphere

## 1. Introduction

Lines of neutral magnesium are present in the Solar spectrum and the corresponding Stark broadening parameters are of interest for their analysis as well as for the diagnostic of Solar plasma. Especially the infrared lines of Mg I have been observed in the Solar spectrum at Kitt Peak and during the Atmos experiment on Spacelab (Brault & Noyes 1983; Chang & Noyes 1983; Farmer & Norton 1989; Jefferies 1991). Other observations of such lines have been described by Deming et al. 1988, 1991; Glenar et al. 1988). Due to the suitability of these lines for the solar atmosphere investigations (see e.g. Van Regemorter & Hoang-Binh 1993) and to the fact that with the increase of the principal quantum number increases the importance of Stark broadening as well, the corresponding Stark widths and shifts are of importance for the structure of the Solar atmosphere research and solar plasma diagnostic (see e.g. Chang et al. 1991; Carlson et al. 1992). Stark broadening data for Mg I lines are also of interest for laboratory plasma research and have been investigated experimentally (Helbig & Kusch 1972; Kusch & Schweicker 1976; Goldbach et al. 1982) and theoretically (Benett & Griem 1971; Griem 1974; Brissaud et al. 1976; Dimitrijević & Konjević 1986; Marasinghe et al. 1986; Regemorter & Hoang Binh 1993; Dimitrijević & Sahal-Bréchet 1995).

By using the semiclassical-perturbation formalism (Sahal-Bréchet 1969a, b), electron-, proton-, and ionized argon-impact line widths and shifts for 99 Mg I multiplets

have been published recently at an electron density of  $10^{15} \text{ cm}^{-3}$ , in order to provide Stark broadening data needed for laboratory plasma research and diagnostic (Dimitrijević & Sahal-Bréchet 1995). Using the same semiclassical approach, we have calculated here electron-, proton-, Mg II-, Si II-, Fe II-, and Ar II-impact line widths and shifts for 267 Mg I multiplets at an electron density of  $10^{11} \text{ cm}^{-3}$ , in order to provide the needed Stark broadening parameters for all important perturbers in the investigation and modelling of Solar plasma. A summary of the formalism is given in Dimitrijević et al. (1991).

## 2. Results and discussion

Energy levels for Mg I lines have been taken from Bashkin & Stoner (1975). For high  $n$  and  $\ell$ , needed atomic energy levels not existing in Bashkin & Stoner, have been calculated by using the theoretical polarization formula of Waller (1926) (see e.g. Regemorter & Hoang Binh 1993 or Edlen 1964) valid for the angular momenta  $\ell$  larger than 3. Oscillator strengths have been calculated by using the method of Bates & Damgaard (1949) and the tables of Oertel & Shomo (1968). For higher levels, the method described by Regemorter et al. (1979) has been used. In order to check the applicability of Bates & Damgaard's (1949) method for the calculation of oscillator strengths within the Coulomb approximation, line width calculations with line strengths from TOP base (the complete package of the opacity project (OP) data with the database management system is usually referred to as TOP base, see Butler et al. (1993) and Cunto et al. (1993)) have been performed at an electron density of  $10^{12} \text{ cm}^{-3}$  and  $T=5000 \text{ K}$ . For the  $3s^2 \ ^1S-3s3p \ ^1P^0$  resonance line the

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<sup>\*</sup>Tables 1-3 are only available in electronic form via anonymous ftp 130.79.128.5

full width with Bates & Damgaard line strengths is  $0.63 \cdot 10^{-6}$  Å and with TOP base line strengths  $0.52 \cdot 10^{-6}$  Å. For  $6s^3S-6p^3P^0$  infrared line the full halfwidth is 0.162 Å in both cases. One can see that for transitions between low levels small differences exist. However, for lines that are of interest for solar infrared astrophysics, corresponding to transitions between high levels, the use of Bates & Damgaard oscillator strengths is sufficient for line widths. In addition to electron-impact full halfwidths and shifts, Stark-broadening parameters due to proton-, Mg II-, Si II- and Fe II- impacts have been calculated in order to provide an as much as possible complete set of Stark broadening data for Solar plasma research. Moreover, Stark-broadening parameters due to Ar II-impacts have been calculated as well, since argon is usually used as a carrier gas for experimental determination of Mg I Stark widths and shifts.

Our results for Stark broadening parameters due to e-, p-, and Ar II-impacts for 267 Mg I multiplets are shown in Table 1 (accessible only in electronic form), and due to Si II-, Fe II-, and Mg II-impacts in Table 2 (accessible only in electronic form) for perturber density of  $10^{11} \text{ cm}^{-3}$  and temperatures  $T=2.500-50.000$  K. Stark broadening parameters due to e-, p-, and Ar II-impacts for perturber densities  $10^{15} \text{ cm}^{-3}-10^{19} \text{ cm}^{-3}$  are shown in Table 3 (accessible only in electronic form). We also specify a parameter  $c$  (Dimitrijević & Sahal-Bréchet 1984), which gives an estimate for the maximum perturber density for which the line may be treated as isolated when it is divided by the corresponding electron-impact full width at half maximum. For each value given in Table 1, the collision volume ( $V$ ) multiplied by the perturber density ( $N$ ) is much less than one and the impact approximation is valid (Sahal-Bréchet 1969a, b). Values for  $NV > 0.5$  are not given and values for  $0.1 < NV \leq 0.5$  are denoted by an asterisk. When the impact approximation is not valid, the ion broadening contribution may be estimated by using quasistatic estimations (Sahal-Bréchet 1991; Griem 1974). The accuracy of the results obtained decreases when broadening by ion interactions becomes important.

Regemorter & Hoang Binh (1993) have discussed recently the important simplification of theory for solar Rydberg lines corresponding to the transitions between nearly degenerate states, where the differences of energy  $E(n\ell, n\ell + 1)$  are of the order of a few  $\text{cm}^{-1}$ . These energy differences are much larger than the width so that the considered lines are isolated. They are however considerably smaller than the energy distances to the other perturbing levels  $n'\ell'$  which may be neglected so that the Stark broadening problem corresponds to a simple three level close coupling system including  $n\ell, n\ell + 1$  atomic energy levels. In order to investigate such approximation, contributions (in angular frequency units) to the Mg I  $5g^1G-6h^1H$  line widths, due to different types of colli-

sions, have been calculated as a function of the number of perturbing levels at an electron density of  $10^{12} \text{ cm}^{-3}$  and temperature of 5000 K and shown in Table 4. One can see that the inelastic collision contribution practically does not vary with the addition of new perturbing levels and that the variation of strong collision contribution is as well of the order of several percents. The elastic collision contribution has however a non negligible dependence on the perturbing level number. Since for solar Rydberg lines the elastic contribution to the width is for an order of magnitude smaller in comparison with other considered contributions, we may conclude that the approximation of Regemorter & Hoang Binh (1993) is very good for Stark widths of the considered lines.

**Table 4.** Contribution of different types of collisions to the line widths (in angular frequency units) as a function of the number of perturbing levels of the initial level ( $i$ ) for Mg I  $5g^1G-6h^1H$  transition at an electron density of  $10^{12} \text{ cm}^{-3}$  for  $T=5000$  K. Perturbing energy levels for the final level are 4f, 5f, 6f, 7f, 6h, 7h, 8h

Perturbing levels $i'$	6g	6g,7f	5g,6g,7f	5g,6g,7g,8g,7f,8f,9f
coll. upper level	0.452(+9)	0.456(+9)	0.456(+9)	0.456(+9)
coll. lower level	0.109(+9)	0.109(+9)	0.109(+9)	0.109(+9)
strong coll.	0.119(+9)	0.135(+9)	0.128(+9)	0.129(+9)
elastic coll.	0.368(+8)	0.590(+8)	0.468(+8)	0.495(+8)

Since the elastic collisions have an important role for the Stark shift determination, for Stark shifts such approximation may become questionable.

In Table 5 our results (denoted as DSB) for Stark full widths for Mg I  $5g^1G-6h^1H$  and  $6g^1G-7h^1H$  transitions have been compared with results of Regemorter & Hoang Binh (1993) (denoted as RHB), with calculation within the Griem's (1968) semiempirical approach of Carlsson et al. (1992) (denoted as GCRS) and with calculations of Chang & Schoenfeld (1991) by using the Lindholm (1941) adiabatic theory (LCS). One can see that our calculations are in good agreement with the width calculations of Regemorter & Hoang Binh (1993) while for shifts the differences due to the omission of perturbing levels with the different principal quantum number (as discussed concerning Table 4) exist. The agreement with adiabatic theory of Lindholm (1941) used by Chang & Schoenfeld (1991), is apparently fortuitous for the width but shifts are largely overestimated, while the Griem's (1968) semiempirical approach, not appropriate for transitions between nearly degenerate states (GCRS), is in agreement with our results

**Table 5.** Comparison of electron-impact width (W-FWHM) and shift ( $d$ ) values calculated according to various approaches, for Mg I  $5g^1G-6h^1H$  and  $6g^1G-7h^1H$  transitions at an electron density of  $10^{12} \text{ cm}^{-3}$  for  $T=5000 \text{ K}$ . The notation is: DSB - present calculations; RHB - Regemorter & Hoang Binh (1993); GCRS - calculated by Carlson et al. (1992) by using the Griem's (1968) semiempirical formula with an effective Gaunt factor  $g = 0.5$ ; LCS - calculated by Chang & Schoenfeld (1991) by using the Lindholm (1941) adiabatic theory

Transition	DSB	RHB	GCRS	LCS
<b>5g - 6h</b>				
$W[\text{Å}]$	0.18	0.16	0.21	0.25
$d[\text{Å}]$	0.0034	0.013	-	0.22
<b>6g - 7h</b>				
$W[\text{Å}]$	2.48	2.67	1.07	2.73
$d[\text{Å}]$	-0.0128	0.052	-	2.36

for Mg I  $5g^1G-6h^1H$  transition but gives more than two times smaller result for  $6g^1G-7h^1H$  transition.

We hope that the comprehensive set of Stark broadening parameters of Mg I lines will enable the better use of Mg I spectral lines for solar plasma research.

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